

Last time: Inverse Kinematics

Lecture 7, ME221, 10/18/2021

Analytic:

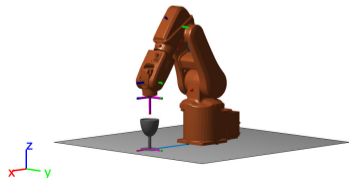
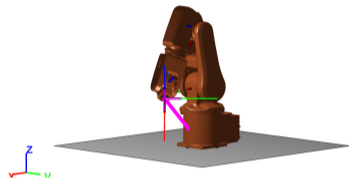
```
% generate a function called 'robotIK'  
generateIKFunction(aik, 'robotIK');  
% find the pose for the target end-effector position  
q = robotIK(Td);
```

Numeric:

```
% BFGSGradientProjection IK object  
ik = inverseKinematics('RigidBodyTree', r);  
  
% find pose with numerical IK  
[q, solnInfo] = ik('link_6', Td, ones(6,1), zeros(6,1));
```

Generalized Numeric:

```
% Generalized IK object  
gik = generalizedInverseKinematics('RigidBodyTree', r);  
gik.ConstraintInputs = {'cartesian', 'position', 'aiming'};  
  
...  
  
% Get IK  
[q, solnInfo] = gik(q0, heightAboveFloor, ...  
                    distanceFromCup, ...  
                    alignWithCup);
```



Today's Agenda

- Time Derivatives
- Angular Velocity Vector / Matrix
- Rigid Body Velocity
- Velocity Transformation Matrix
- Homogeneous Transformation Derivatives

Time Derivatives and Coordinate Frames

General case: The time derivative of a vector depends on *which* coordinate frame we want to take the derivative in:

G-derivative:

(Derivative in the Global Frame)

$$\frac{Gd}{dt} \mathbf{r}_p$$

B-derivative:

Derivative in the Body Frame

$$\frac{Bd}{dt} \mathbf{r}_p$$

Time Derivatives and Coordinate Frames

If vector \mathbf{r}_p is represented in the same coordinate frame as the derivative, the frame's *unit vectors* are constant:

G-derivative:

(Derivative in the Global Frame)

$$\begin{aligned}\frac{{}^G d}{dt} {}^G \mathbf{r}_p &= {}^G \dot{\mathbf{r}}_p = {}^G \mathbf{v}_p \\ &= \dot{X} \hat{I} + \dot{Y} \hat{J} + \dot{Z} \hat{K}\end{aligned}$$

B-derivative:

Derivative in the Body Frame

$$\begin{aligned}\frac{{}^B d}{dt} {}^B \mathbf{r}_p &= {}^B \dot{\mathbf{r}}_p = {}^B \mathbf{v}_p \\ &= \dot{x} \hat{i} + \dot{y} \hat{j} + \dot{z} \hat{k}\end{aligned}$$

Time Derivatives and Coordinate Frames

We can also find the **G-derivative** of ${}^B\mathbf{r}_p$:

G-derivative:

(Derivative in the Global Frame)

$${}^G\mathbf{v}_p = \frac{{}^G d}{dt} {}^B\mathbf{r}_p = {}^B\dot{\mathbf{r}}_p + {}^B_G\boldsymbol{\omega}_B \times {}^B\mathbf{r}_p$$

where ${}^B_G\boldsymbol{\omega}_B$ is the angular velocity of B relative to G.

How to derive?

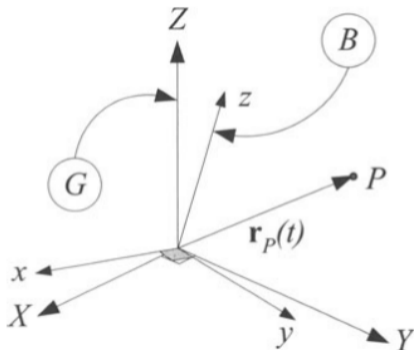


Figure: 7.3 - Moving point P in rotating frame B

Position in the body frame:

$${}^B \mathbf{r}_p = x(t)\hat{i} + y(t)\hat{j} + z(t)\hat{k}$$

Take time derivative:

$$\begin{aligned} \frac{{}^G d}{dt} {}^B \mathbf{r}_p &= \frac{{}^G d}{dt} (x\hat{i} + y\hat{j} + z\hat{k}) \\ &= \dot{x}\hat{i} + \dot{y}\hat{j} + \dot{z}\hat{k} + x \frac{{}^G d\hat{i}}{dt} + y \frac{{}^G d\hat{j}}{dt} + z \frac{{}^G d\hat{k}}{dt} \end{aligned}$$

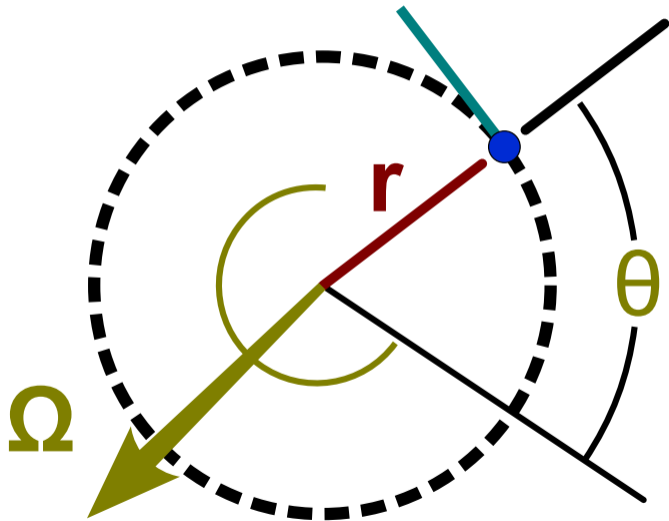
Time derivative of unit vectors:

$$\frac{{}^G d\hat{i}}{dt} = {}^B \boldsymbol{\omega}_B \times \hat{i}, \quad \frac{{}^G d\hat{j}}{dt} = {}^B \boldsymbol{\omega}_B \times \hat{j}, \quad \frac{{}^G d\hat{k}}{dt} = {}^B \boldsymbol{\omega}_B \times \hat{k}$$

Therefore:

$${}^B \mathbf{v}_p = {}^B \dot{\mathbf{r}}_p + {}^B \boldsymbol{\omega}_B \times {}^B \mathbf{r}_p$$

$$\mathbf{v} = \boldsymbol{\Omega} \times \mathbf{r}$$



Time Derivatives and Coordinate Frames

G-derivative of ${}^B\mathbf{r}_p$ and the **B-derivative** of ${}^G\mathbf{r}_p$:

G-derivative:

(Derivative in the Global Frame)

$$\begin{aligned} {}^G\mathbf{v}_p &= \frac{{}^G d}{dt} {}^B\mathbf{r}_p \\ &= {}^B\dot{\mathbf{r}}_p + {}^B\boldsymbol{\omega}_B \times {}^B\mathbf{r}_p \end{aligned}$$

where ${}^B\boldsymbol{\omega}_B$ is the angular velocity of B relative to G.

B-derivative:

(Derivative in the Body Frame)

$$\begin{aligned} {}^G\mathbf{v}_p &= \frac{{}^B d}{dt} {}^G\mathbf{r}_p \\ &= {}^G\dot{\mathbf{r}}_p + {}^G\boldsymbol{\omega}_G \times {}^G\mathbf{r}_p \end{aligned}$$

where ${}^G\boldsymbol{\omega}_G$ is the angular velocity of G relative to B.

Angular Velocity Vector / Matrix

Specific case: Consider a rigid body rotating about a fixed point. In this case ${}^B\mathbf{r}_p$ represents a point on the rigid body (e.g., center of mass), thus is constant.

$${}^G\mathbf{r}_p(t) = {}^G R_B(t) {}^B\mathbf{r}_p$$

What is the velocity of point p ?

$$\begin{aligned} {}^G\mathbf{v}_p(t) &= {}^G\dot{\mathbf{r}}_p(t) = {}^G\dot{R}_B(t) {}^B\mathbf{r}_p \\ &= {}_G\boldsymbol{\omega}_B \times {}^G\mathbf{r}_p(t) \\ &= {}_G\tilde{\boldsymbol{\omega}}_B {}^G\mathbf{r}_p(t) \end{aligned}$$

where the skew symmetric matrix:

$${}_G\tilde{\boldsymbol{\omega}}_B = {}^G\dot{R}_B(t) {}^G R_B^T(t) = \dot{\phi}\tilde{\mathbf{u}}$$

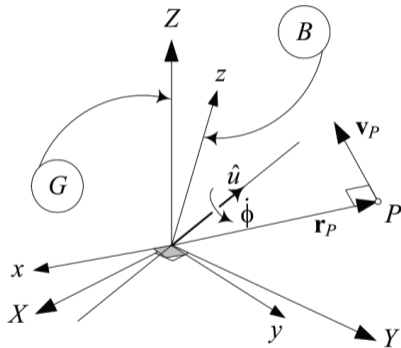


Figure: 7.1 - Rotating body frame.

$${}^G \mathbf{v}_p(t) = {}^G \dot{\mathbf{r}}_p(t) = {}^G \dot{R}_B(t) {}^B \mathbf{r}_p$$

But we can write:

$${}^B \mathbf{r}_p(t) = {}^G R_B^T(t) {}^G \mathbf{r}_p$$

so

$${}^G \mathbf{v}_p(t) = {}^G \dot{\mathbf{r}}_p(t) = \underbrace{{}^G \dot{R}_B(t) {}^G R_B^T(t)}_{{}^G \tilde{\omega}_B} {}^G \mathbf{r}_p(t)$$

Why should ${}^G \tilde{\omega}_B$ be skew symmetric?

Start from orthogonality conditions:

$$\begin{array}{ll}
 {}^G R_B {}^G R_B^T = I & {}^G R_B^T {}^G R_B = I \\
 \frac{d}{dt} ({}^G R_B {}^G R_B^T = I) & \frac{d}{dt} ({}^G R_B^T {}^G R_B = I) \\
 {}^G \dot{R}_B {}^G R_B^T + {}^G R_B {}^G \dot{R}_B^T = 0 & {}^G \dot{R}_B^T {}^G R_B + {}^G R_B^T {}^G \dot{R}_B = 0 \\
 {}^G \dot{R}_B {}^G R_B^T = - {}^G R_B {}^G \dot{R}_B^T & {}^G \dot{R}_B^T {}^G R_B = - {}^G R_B^T {}^G \dot{R}_B \\
 ({}^G R_B {}^G \dot{R}_B^T)^T = - {}^G R_B {}^G \dot{R}_B^T & ({}^G R_B^T {}^G \dot{R}_B)^T = - {}^G R_B^T {}^G \dot{R}_B
 \end{array}$$

Note that:

$$B^T A = (A^T B)^T \quad \text{and} \quad A^T = -A \quad \text{for skew symmetric matrix}$$

The group of all skew-symmetric 3×3 matrices forms the Lie algebra $\mathfrak{so}(3)$ of the Lie group $\overline{SO(3)}$ (special rotation group). They represent *infinitesimal generators* of rotations, e.g., **angular velocity!**

$${}^G \dot{\mathbf{r}}_p(t) = {}^G \dot{R}_B(t) {}^B \mathbf{r}_p$$

Left multiply rotation matrix:

$$\underbrace{{}^G R_B^T(t) {}^G \dot{\mathbf{r}}_p(t)}_{{}^B \dot{\mathbf{r}}_p} = \underbrace{{}^G R_B^T(t) {}^G \dot{R}_B(t)}_{{}^B \tilde{\omega}_B} {}^B \mathbf{r}_p$$

Substitute ${}^G \mathbf{r}_p$ on RHS:

$${}^G \dot{\mathbf{r}}_p(t) = \underbrace{{}^G \dot{R}_B(t) {}^G R_B^T(t)}_{{}^G \tilde{\omega}_B} {}^G \mathbf{r}_p$$

${}^B \tilde{\omega}_B$: instantaneous angular velocity of B relative to G as seen from B.

${}^G \tilde{\omega}_B$: instantaneous angular velocity of B relative to G as seen from G.

Rigid Body Velocity

Now consider a rigid body rotating and translating (freely moving) relative to the Global frame.

$${}^G \mathbf{r}_p(t) = {}^G R_B(t) {}^B \mathbf{r}_p + {}^G \mathbf{d}_B$$

What is the velocity of point p in the Global frame?

$$\begin{aligned} {}^G \mathbf{v}_p(t) &= {}^G \dot{\mathbf{r}}_p(t) = {}^G \dot{R}_B(t) {}^B \mathbf{r}_p + {}^G \dot{\mathbf{d}}_B \\ &= {}^G \tilde{\omega}_B ({}^G \mathbf{r}_p - {}^G \mathbf{d}_B) + {}^G \dot{\mathbf{d}}_B \end{aligned}$$

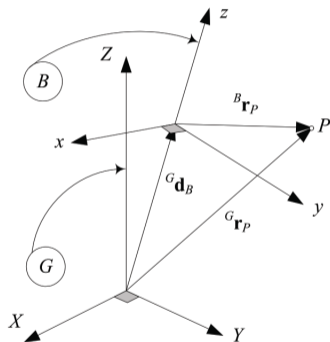


Figure: 7.5 - Rotating and translating body frame.

Geometric Interpretation of Rigid Body Velocity

$${}^G \mathbf{v}_p = {}_G \tilde{\omega}_B ({}^G \mathbf{r}_p - {}^G \mathbf{d}_B) + {}^G \dot{\mathbf{d}}_B$$

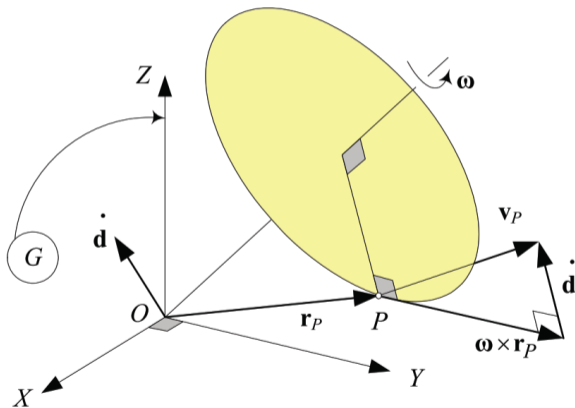


Figure: 7.6 - Rigid body velocity.

Velocity Transformation Matrix

The Velocity Transformation Matrix is a more convenient representation of rigid body velocity by utilizing a matrix operator

$${}^G \mathbf{v}_p(t) = {}^G V_B {}^G \mathbf{r}_p$$

Note that we need to represent the point in the Global frame

$$\begin{aligned} {}^G V_B &= {}^G \dot{T}_B {}^G T_B^{-1} \\ &= \begin{bmatrix} {}^G \dot{R}_B & {}^G R_B^T & {}^G \dot{\mathbf{d}}_B - {}^G \dot{R}_B {}^G R_B^T {}^G \mathbf{d}_B \\ 0 & 0 & 0 \end{bmatrix} \\ &= \begin{bmatrix} {}^G \tilde{\omega}_B & {}^G \dot{\mathbf{d}}_B - {}^G \tilde{\omega}_B {}^G \mathbf{d}_B \\ 0 & 0 & 0 \end{bmatrix} \\ &= \begin{bmatrix} {}^G \tilde{\omega}_B & {}^G \mathbf{v}_B \\ 0 & 0 \end{bmatrix} \quad (\text{same as } \tilde{S}!) \end{aligned}$$

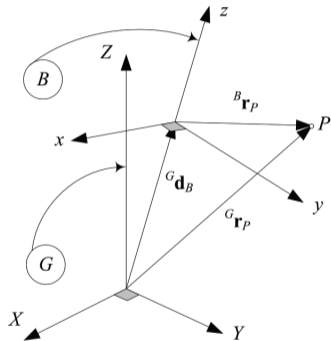


Figure: 7.5 - Rotating and translating body frame.

Velocity Transformation Matrix

The Velocity Transformation Matrix is a matrix operator that provides the global velocity of any point attached to the B frame.

$${}^G \mathbf{v}_p(t) = {}^G V_B {}^G \mathbf{r}_p$$

We can also represent as a velocity transformation vector, also called a *twist*:

$${}^G V_B = \begin{bmatrix} {}^G \tilde{\omega}_B & {}^G \mathbf{v}_B \\ 0 & 0 \end{bmatrix}$$
$${}^G \mathbf{t}_B = \begin{bmatrix} {}^G \mathbf{v}_B \\ {}^G \boldsymbol{\omega}_B \end{bmatrix}$$

The group of all twist matrices form the Lie algebra $\mathfrak{se}(3)$ of the Lie group $SE(3)$ (special Euclidean group). They represent *infinitesimal generators* of screw motions, e.g., **angular and translational velocity!**

Derivative of a Homogeneous Transformation Matrix

The velocity transformation matrix is related to the derivative of a homogeneous transformation matrix (e.g., can be used for successive link transformation for velocity forward kinematics)

$${}^G\dot{T}_B = {}^G V_B {}^G T_B$$

$${}^{i-1}\dot{T}_i = {}^{i-1} V_i {}^{i-1} T_i$$

$${}^{i-1} V_i = \underbrace{{}^{i-1}\dot{T}_i {}^{i-1} T_i^{-1}}_{\text{known for each joint!}}$$

Next time we will apply this to robot manipulators for velocity (forward/inverse) kinematics.

Informal Instructor Feedback Exercise

- Do you feel that you are learning?
- What do you like about the course? What is helping you learn?
- What do you not like about the course? What is inhibiting your learning?
- What changes would you implement if you were the instructor?
 - The homeworks are too hard (or too easy).
 - Lectures are boring. I would make them more exciting.
 - I would implement slides only instead of handwritten notes.
 - etc.