

Quantum Un-Demolition: Undoing quantum measurement by erasing information

Alexander Korotkov¹ and Andrew Jordan²

¹*University of California, Riverside*

²*University of Rochester*

Outline:

- The problem
- QUD for a charge qubit
- QUD for a phase qubit
- General theory of QUD

Support:



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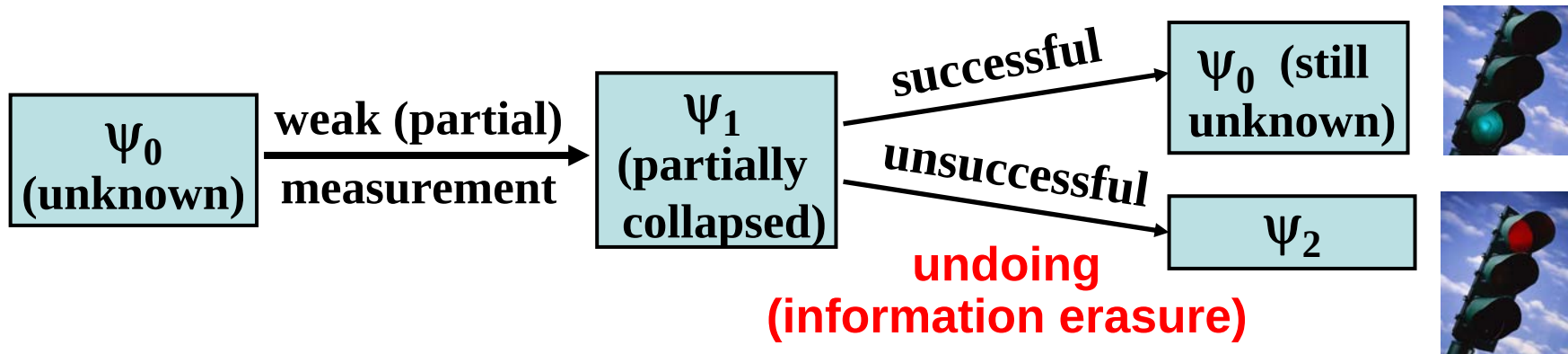
The problem

It is impossible to undo “orthodox” quantum measurement (for an unknown initial state)

Is it possible to undo weak (partial) quantum measurement?
(To restore a “precious” qubit accidentally measured)

Yes! (but with a finite probability)

If undoing is successful, an unknown state is **fully** restored

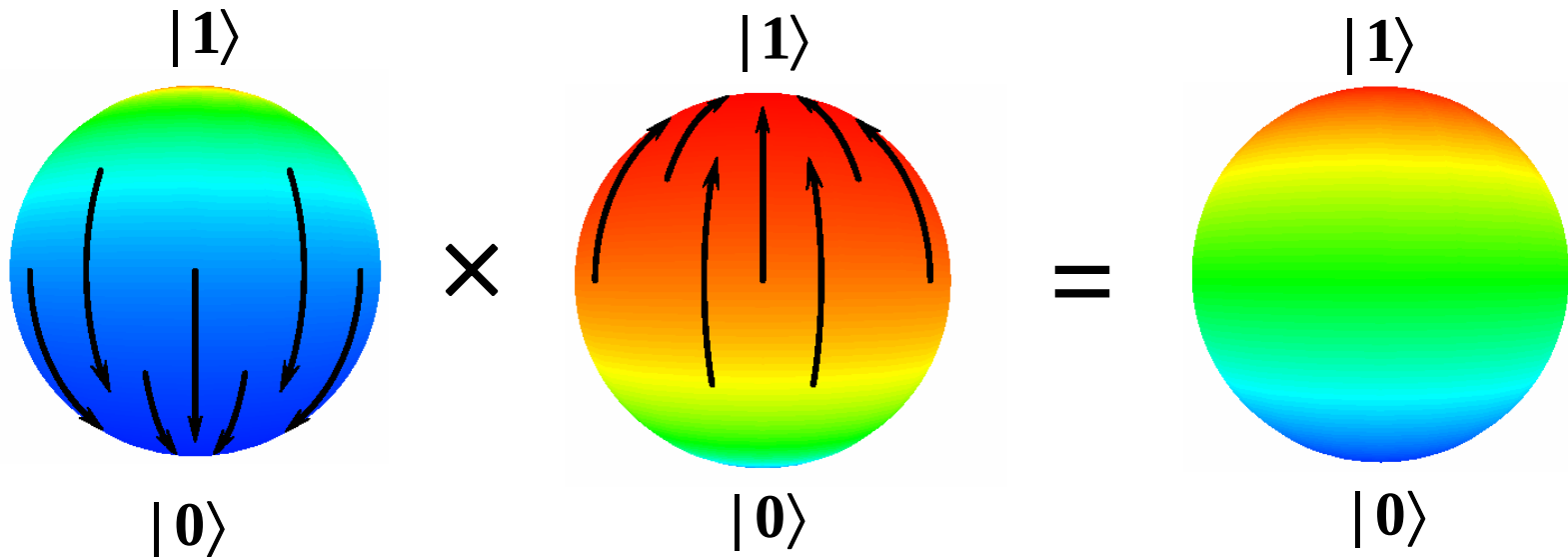


“Quantum Un-Demolition (QUD) measurement”

Measurement undoing

Evolution due to partial (weak, continuous, etc.) measurement is **non-unitary** (though coherent if detector is good!), therefore it is impossible to undo it by Hamiltonian dynamics.

How to undo? One more measurement!

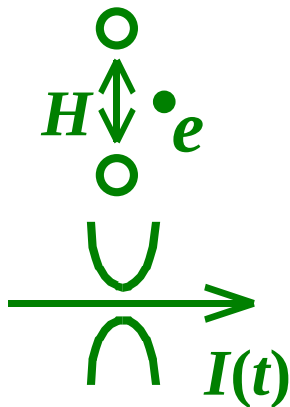


(similar to Koashi-Ueda, PRL-1999)

(Figure partially adopted from A. Jordan, A. Korotkov, and M. Büttiker, PRL-2006)



First example: DQD qubit with no tunneling, measured by QPC



$$\hat{H}_{QB} = (\varepsilon/2)(c_1^\dagger c_1 - c_2^\dagger c_2) + H(c_1^\dagger c_2 + c_2^\dagger c_1)$$

Assume “frozen” qubit: $\varepsilon = H = 0$

Bayesian evolution due to measurement (Korotkov-1998)

- 1) Diagonal matrix elements of the density matrix evolve according to the classical Bayes rule
- 2) Non-diagonal matrix elements evolve so that the degree of purity $\rho_{ij}/[\rho_{ii} \rho_{jj}]^{1/2}$ is conserved

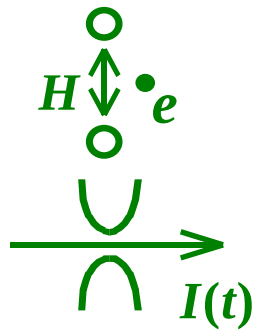
$$\rho_{11}(\tau) = \frac{\rho_{11}(0) \exp[-(\bar{I} - I_1)^2 / 2D]}{\rho_{11}(0) \exp[-(\bar{I} - I_1)^2 / 2D] + \rho_{22}(0) \exp[-(\bar{I} - I_2)^2 / 2D]}$$

$$\frac{\rho_{12}(\tau)}{[\rho_{12}(\tau) \rho_{22}(\tau)]^{1/2}} = \frac{\rho_{12}(0)}{[\rho_{12}(0) \rho_{22}(0)]^{1/2}}, \quad \rho_{22}(\tau) = 1 - \rho_{11}(\tau)$$

where $\bar{I} = \frac{1}{\tau} \int_0^\tau I(t) dt$



Graphical representation of the evolution

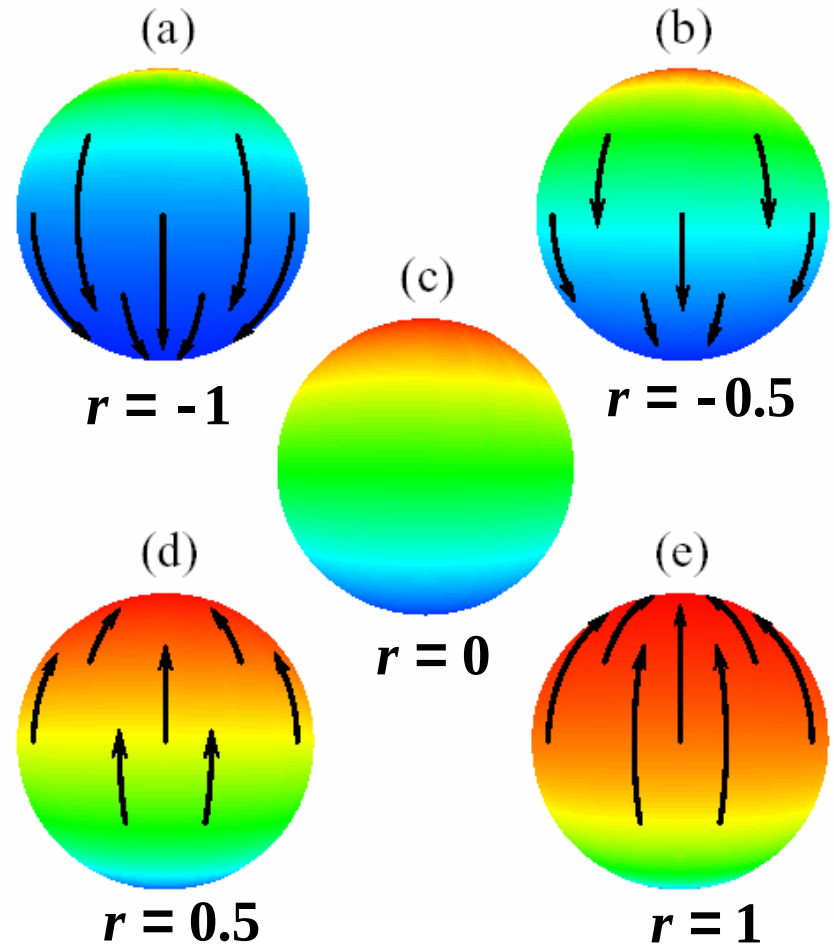


$$\frac{\rho_{11}(t)}{\rho_{22}(t)} = \frac{\rho_{11}(0)}{\rho_{22}(0)} \exp[2r(t)]$$

$$\frac{\rho_{12}(t)}{\sqrt{\rho_{11}(t)\rho_{22}(t)}} = \text{const}$$

where measurement result $r(t)$ is

$$r(t) = \frac{\Delta I}{S_I} \left[\int_0^t I(t') dt' - I_0 t \right]$$



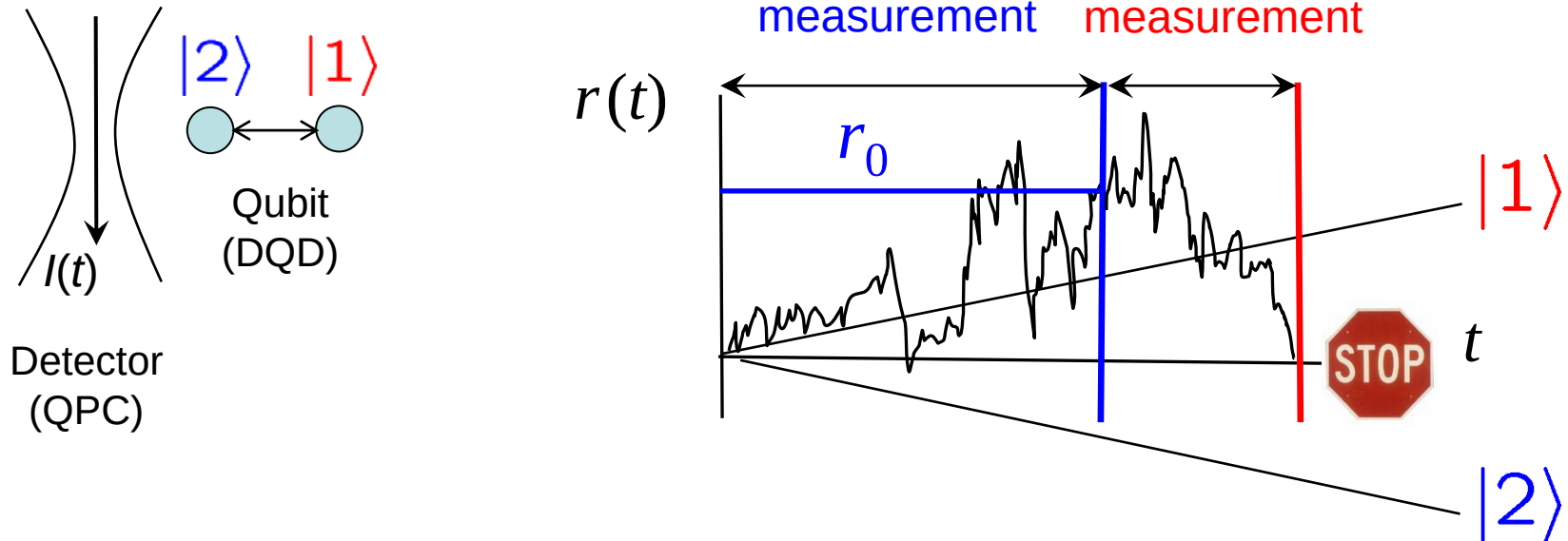
Jordan-Korotkov-Büttiker, PRL-06

If $r = 0$, then no information and no evolution!



Measurement undoing for DQD-QPC system

Jordan and Korotkov, PRL-2006



**Simple strategy: continue measuring until $r(t)$ becomes zero!
Then any unknown initial state is fully restored.**

(same for an entangled qubit)

It may happen though that $r = 0$ never happens;
then undoing procedure is unsuccessful.



Probability of success

Trick: since non-diagonal matrix elements are not directly involved, we can analyze classical probabilities (as if qubit is in some certain, but unknown state); then simple diffusion with drift

Results:

**Probability of
successful undoing**

$$P_s = \frac{e^{-|r_0|}}{e^{|r_0|}\rho_{11}(0) + e^{-|r_0|}\rho_{22}(0)}$$

where r_0 is the result of the measurement to be undone, and $\rho(0)$ is our knowledge about an unknown initial state; in case of an entangled qubit $\rho(0)$ is traced over other qubits

Average time to wait

$$T_{\text{undo}} = T_m |r_0| \quad \text{where} \quad T_m = 2S_I / (\Delta I)^2$$

("measurement time")

**Averaged probability
of success (over result r_0)**

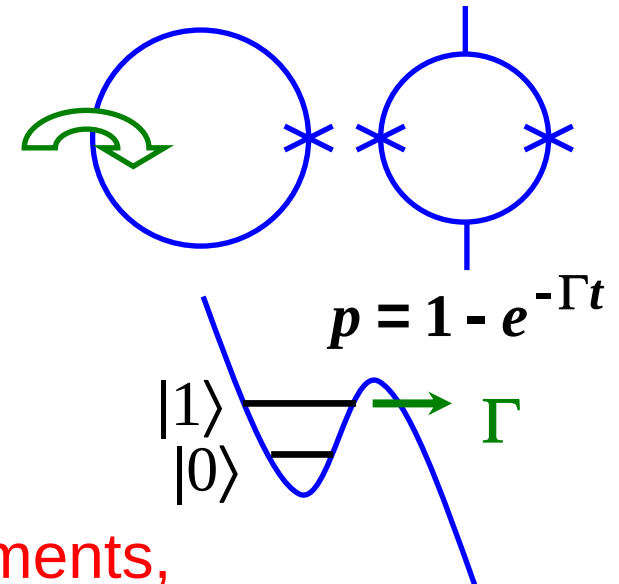
$$P_{\text{av}} = 1 - \text{erf}[\sqrt{t / 2T_m}]$$

(does not depend on initial state!)



Second example: Undoing partial measurement of a phase qubit

- 1) Start with an unknown state
- 2) Partial measurement of strength p
- 3) π -pulse (exchange $|0\rangle \leftrightarrow |1\rangle$)
- 4) One more measurement with the **same strength p**
- 5) π -pulse



If no tunneling for both measurements,
then initial state is fully restored!

N. Katz et al.,
Science-2006

$$\alpha |0\rangle + \beta |1\rangle \rightarrow \frac{\alpha |0\rangle + e^{i\phi} \beta e^{-\Gamma t/2} |1\rangle}{\text{Norm}} \rightarrow$$

$$\frac{e^{i\phi} \alpha e^{-\Gamma t/2} |0\rangle + e^{i\phi} \beta e^{-\Gamma t/2} |1\rangle}{\text{Norm}} = e^{i\phi} (\alpha |0\rangle + \beta |1\rangle)$$



Probability of successful measurement undoing for phase qubit

Success probability if no tunneling during first measurement:

$$P_s = \frac{e^{-\Gamma t}}{\rho_{00}(0) + e^{-\Gamma t} \rho_{11}(0)} = \frac{1 - p}{\rho_{00}(0) + (1 - p) \rho_{11}(0)}$$

where $\rho(0)$ is the density matrix of the initial state (either averaged unknown state or an entangled state traced over all other qubits)

Total (averaged) success probability: $P_{av} = 1 - p$

For measurement strength p increasing to 1, success probability decreases to zero (orthodox collapse), but still exact undoing

**Such an experiment is only slightly more difficult than recent experiment on partial collapse (N. Katz et al., 2006).
Can be realized experimentally pretty soon!!!**



General theory of quantum measurement undoing

Measurement operator M_r : $\rho \rightarrow \frac{M_r \rho M_r^\dagger}{\text{Tr}(M_r \rho M_r^\dagger)}$ (POVM formalism)

Undoing measurement operator: $C \times M_r^{-1}$ (to satisfy completeness, eigenvalues cannot be >1)

$$\max(C) = \min_i \sqrt{p_i}, \quad p_i = \text{Tr}(M_r^\dagger M_r |i\rangle \langle i|)$$

p_i – probability of the measurement result r for initial state $|i\rangle$

Probability of success:
$$P_s \leq \frac{\min_i p_i}{\sum_i p_i \rho_{ii}(0)} = \frac{\min P_r}{P_r(\rho(0))}$$

$P_r(\rho(0))$ – probability of result r for initial state $\rho(0)$,

$\min(P_r)$ – probability of result r minimized over all possible initial states

Averaged (over r) probability of success:
$$P_{av} \leq \sum_r \min P_r$$

(similar to Koashi-Ueda, PRL, 1999)



Comparison of the general bound for undoing success with examples

General bound:
$$P_S \leq \frac{\min P_r}{P_r(\rho(0))}$$

First example (DQD+QPC)
$$P_S \leq \frac{\min(p_1, p_2)}{p_1 \rho_{11}(0) + p_2 \rho_{22}(0)}$$

where
$$p_i = (\pi S_I / t)^{-1/2} \exp[-(\bar{I} - I_i)^2 t / S_I] d\bar{I}$$

Coincides with the pervious result, so the upper bound is reached,
therefore undoing strategy is optimal

Second example (phase qubit) Probabilities of no-tunneling are 1 and $\exp(-\Gamma t)=1-p$

$$P_S \leq \frac{1-p}{\rho_{00}(0) + (1-p)\rho_{11}(0)}$$

Again same as before, so measurement
undoing for phase qubit is also optimal



Quantum erasers in optics

Quantum eraser proposal by Scully and Drühl, PRA (1982)

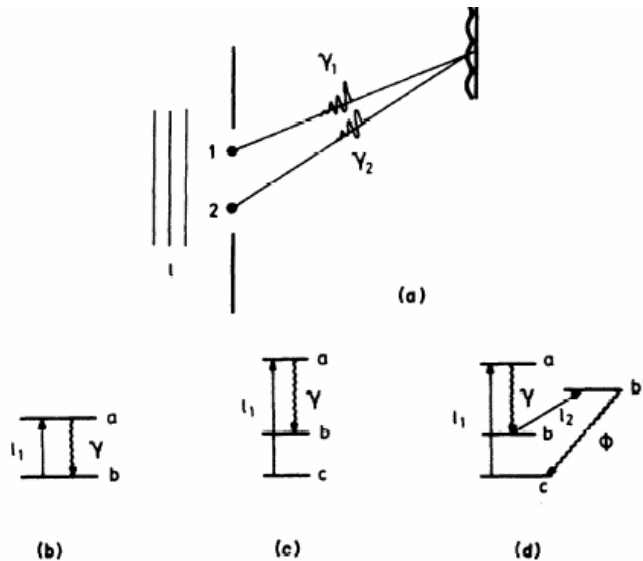


FIG. 1. (a) Figure depicting light impinging from left on atoms at sites 1 and 2. Scattered photons γ_1 and γ_2 produce interference pattern on screen. (b) Two-level atoms excited by laser pulse l_1 , and emit γ photons in $a \rightarrow b$ transition. (c) Three-level atoms excited by pulse l_1 from $c \rightarrow a$ and emit photons in $a \rightarrow b$ transition. (d) Four-level system excited by pulse l_1 from $c \rightarrow a$ followed by emission of γ photons in $a \rightarrow b$ transition. Second pulse l_2 takes atoms from $b \rightarrow b'$. Decay from $b' \rightarrow c$ results in emission of ϕ photons.

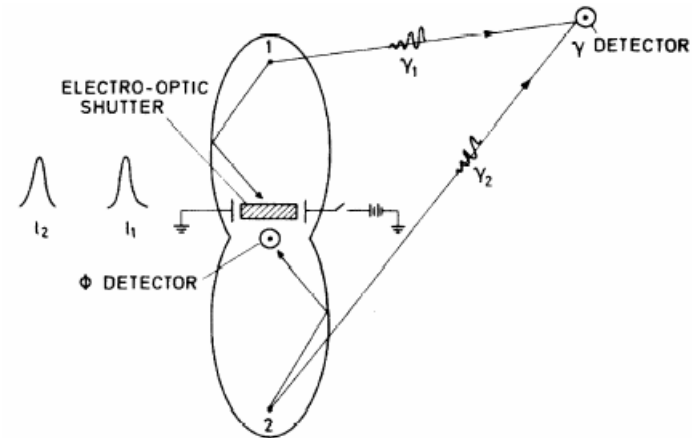


FIG. 2. Laser pulses l_1 and l_2 incident on atoms at sites 1 and 2. Scattered photons γ_1 and γ_2 result from $a \rightarrow b$ transition. Decay of atoms from $b' \rightarrow c$ results in ϕ photon emission. Elliptical cavities reflect ϕ photons onto common photodetector. Electro-optic shutter transmits ϕ photons only when switch is open. Choice of switch position determines whether we emphasize particle or wave nature of γ photons.

Interference fringes restored for two-detector correlations (since “which-path” information is erased)

Our idea of measurement undoing is quite different:
we really extract information and then erase it

Conclusions

- Partial (incomplete, weak, etc.) quantum measurement can be undone, though with a finite probability P_s , which decreases with increasing strength of the measurement ($P_s = 0$ for orthodox case)
- Though somewhat similar to the quantum eraser, undoing idea is actually quite different: “Quantum Un-Demolition” (QUD)
- **Quantum Un-Demolition for a phase qubit can be realized now**, experiment with a charge qubit will hopefully be possible soon (difficulty to use SET: need an ideal quantum detector)

