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Continuous quantum measurement of solid-state systems

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Outline:

- What is "inside" wavefunction collapse (quantum Bayesian formalism)
- Experimental consequences (predictions)
- Recent experiments on partial collapse and "wavefunction uncollapsing"

Support:



Motivation and relevance

- quantum measurement: necessary step in any QIP; also fundamentally interesting (still controversial after 80 years(!), becoming possible to check experimentally)
- solid-state: big potential impact on practical electronics

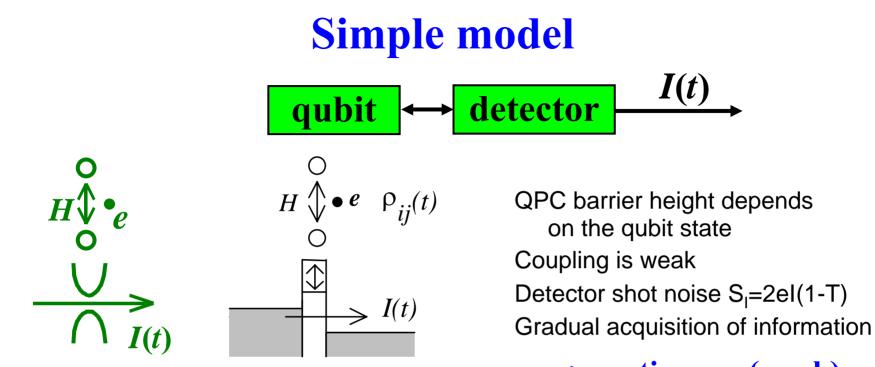
Meanings of "quantum measurement"

- any measurement of a quantum system
- "orthodox" projective measurement (not easy!)
- <u>beyond-orthodox measurement</u> (generalized, POVM, partial, continuous, weak, etc.)

Problem: How does collapse develop in time? (What is "inside" collapse?)

In solid state collapse is typically gradual \Rightarrow important to know





Double-quantum-dot (DQD) as a qubit and quantum point contact (QPC) as a detector

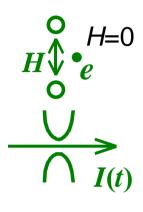
⇒ continuous (weak) measurement

What is the evolution of the DQD qubit state ρ_{ii} ?

- Orthodox projection? No. (Weak coupling, should be gradual evolution.)
- Decoherence? No. (Surprisingly, single qubit retains purity.)
- Evolution is gradual, coherent, and information-related! (A.K., 1998)



Quantum Bayesian formalism for DQD-QPC (qubit-detector) system



Qubit evolution due to continuous measurement:

- 1) Diagonal matrix elements of the qubit density matrix evolve as classical probabilities (i.e. according to the classical Bayes rule)
- 2) Non-diagonal matrix elements evolve so that the degree of purity $\rho_{ij}/[\rho_{ii}\rho_{jj}]^{1/2}$ is conserved

(A.K., 1998)

Bayes rule: $P(A_i | R) = \frac{P(A_i)P(R | A_i)}{\sum_k P(A_k)P(R | A_k)}$

So simple because:

QPC happens to be an ideal detector
 no Hamiltonian evolution of the qubit

Cannot be derived from Schrödinger equation only

Similar formalisms developed earlier. Key words: Imprecise, weak, selective, or conditional measurements, POVM, Quantum trajectories, Quantum jumps, Restricted path integral, etc.

Names: Davies, Kraus, Holevo, Mensky, Caves, Gardiner, Carmichael, Plenio, Knight, Walls, Gisin, Percival, Milburn, Wiseman, Habib, etc. (very incomplete list)



Bayesian formalism for a single qubit $\hat{H}_{QB} = \frac{\varepsilon}{2} (c_1^{\dagger} c_1 - c_2^{\dagger} c_2) + H(c_1^{\dagger} c_2 + c_2^{\dagger} c_1)$ I(t) $|1\rangle \rightarrow I_1, |2\rangle \rightarrow I_2, \Delta I = I_1 - I_2, I_0 = (I_1 + I_2)/2$ S_I – detector noise $\rho_{11} = -\rho_{22} = -2(H/\hbar) \operatorname{Im} \rho_{12} + \rho_{11}\rho_{22}(2\Delta I/S_I)[\underline{I(t)} - I_0]$ $\dot{\rho}_{12} = i(\varepsilon/\hbar)\rho_{12} + i(H/\hbar)(\rho_{11} - \rho_{22}) + \rho_{12}(\rho_{11} - \rho_{22})(\Delta I/S_I)[\underline{I(t)} - I_0] - \gamma\rho_{12}$ (A.K., 1998) $\gamma = \Gamma - (\Delta I)^2 / 4S_I$, Γ – ensemble decoherence $\eta = 1 - \gamma / \Gamma = (\Delta I)^2 / 4S_I \Gamma$ - detector ideality (efficiency), $\eta \le 100\%$ Ideal detector (η =1, as QPC) does not decohere a qubit, then random evolution of qubit wavefunction can be monitored $d\rho_{11}/dt = -d\rho_{22}/dt = -2(H/\hbar) \operatorname{Im} \rho_{12}$ Averaging over result I(t) leads to conventional master equation: $d\rho_{12}/dt = i(\varepsilon/\hbar)\rho_{12} + i(H/\hbar)(\rho_{11} - \rho_{22}) - \Gamma\rho_{12}$ ensemble decoherence always faster Fundamental limit: $\Gamma \ge (\Delta I)^2/4S_I$ than rate of information acquisition Simple generalizations to entangled qubits and other systems

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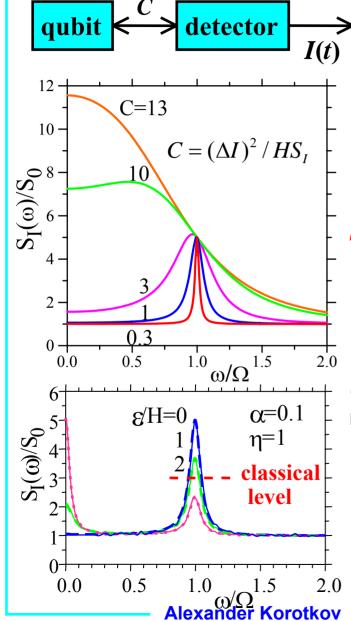
Experimental predictions and proposals from Bayesian formalism

- Direct experimental verification (1998)
- Measured spectral density of Rabi oscillations (1999, 2000, 2002)
- Bell-type correlation experiment (2000)
- Quantum feedback control of a qubit (2001)
- Entanglement by measurement (2002)
- Measurement by a quadratic detector (2003)
- Simple quantum feedback of a qubit (2004)
- Squeezing of a nanomechanical resonator (2004)
- Violation of Leggett-Garg inequality (2005)
- Partial collapse of a phase qubit (2005)
- Undoing of a weak measurement (2006)



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Measured spectrum of Rabi oscillations



What is the spectral density $S_{l}(\omega)$ of detector current?

Assume classical output, eV » $\hbar\Omega$ A $\varepsilon = 0$, $\Gamma = \eta^{-1} (\Delta I)^2 / 4S_0$

 $S_{I}(\omega) = S_{0} + \frac{\Omega^{2} (\Delta I)^{2} \Gamma}{(\omega^{2} - \Omega^{2})^{2} + \Gamma^{2} \omega^{2}}$

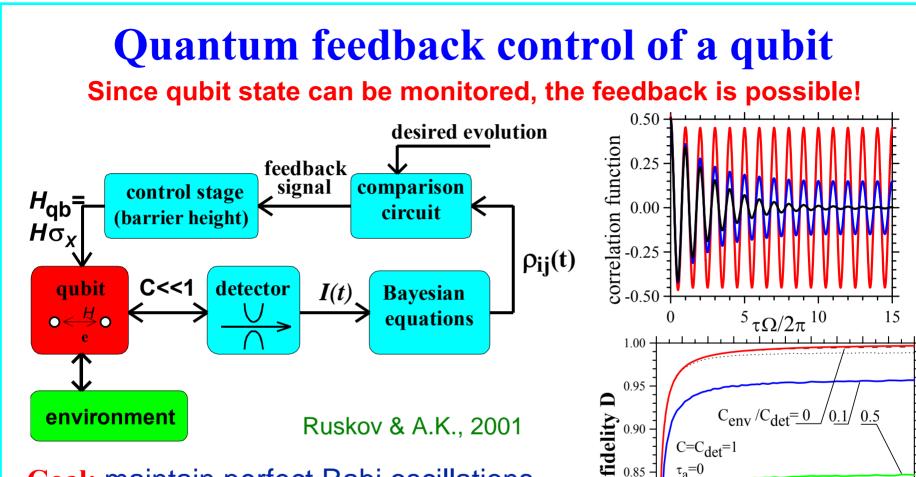
Spectral peak can be seen, but peak-to-pedestal ratio $\leq 4\eta \leq 4$

(result can be obtained using various methods, not only Bayesian method)

A.K., LT'99 A.K.-Averin, 2000 A.K., 2000 Averin, 2000 Goan-Milburn, 2001 Makhlin et al., 2001 Balatsky-Martin, 2001 Ruskov-A.K., 2002 Mozyrsky et al., 2002 Balatsky et al., 2002 Bulaevskii et al., 2002 Shnirman et al., 2003

Spectral peak \Rightarrow Rabi oscillations persist forever if measured (but phase fluctuates)





Goal: maintain perfect Rabi oscillations

Idea: monitor the Rabi phase ϕ by continuous measurement and apply feedback control of the qubit barrier height, $\Delta H_{FB}/H = -F \times \Delta \phi$

First experimental quantum feedback in optics: JM Geremia et al., Science 304, 270 (2004).

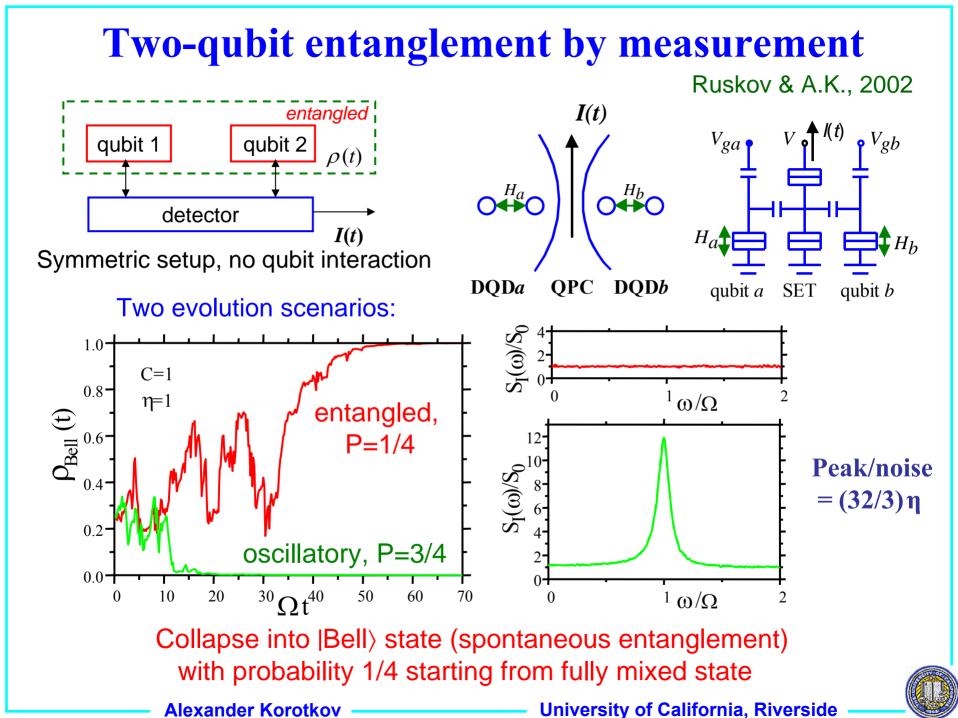
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For ideal detector and wide bandwidth, fidelity can approach 100%

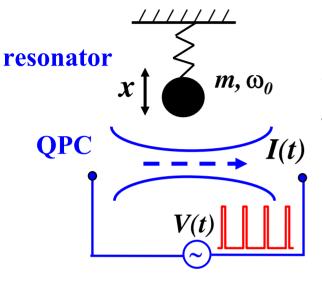
 $= \exp(-C/32F)$

F (feedback strength)

0.80



QND squeezing of a nanomechanical resonator



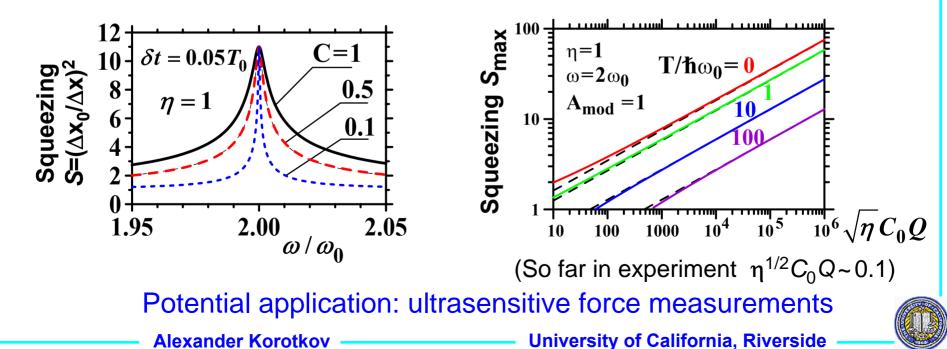
Experimental status:

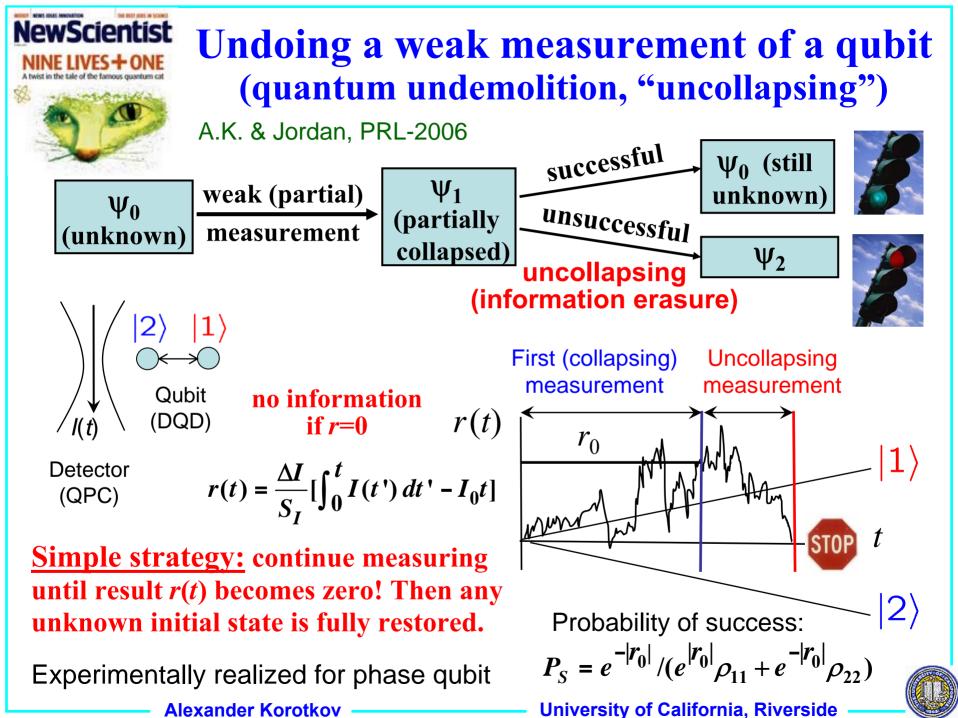
 $ω_0/2π \sim 1 \text{ GHz}$ ($\hbar ω_0 \sim 80 \text{ mK}$), Roukes' group, 2003 $\Delta x/\Delta x_0 \sim 5 \text{ [SQL } \Delta x_0 = (\hbar/2mω_0)^{1/2}\text{]}$, Schwab's group, 2004

$$S_{\max} = \frac{3}{4} \left[\frac{\sqrt{\eta} C_0 Q}{\coth(\hbar \omega_0 / 2T)} \right]^{1/3}$$

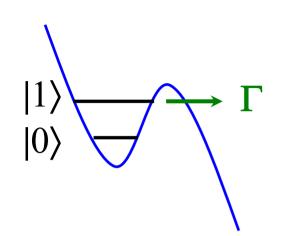
 $m{C}_0$ – coupling with detector, η – detector efficiency, T – temperature, Q – resonator Q-factor

Ruskov, Schwab, Korotkov, PRB-2005





Partial collapse of a superconducting phase qubit



Katz, Ansmann, Bialczak, Lucero, McDermott, Neeley, Steffen, Weig, Cleland, Martinis, Korotkov, *Science-06*

How does a wavefunction evolve in time before tunneling event?

(What happens when nothing happens?)

Qubit "ages" in contrast to a radioactive atom!

Main idea:

$$\psi = \alpha | 0 \rangle + \beta | 1 \rangle \rightarrow \psi(t) = \begin{cases} |out\rangle, \text{ if tunneled} \\ \frac{\alpha | 0 \rangle + \beta e^{-\Gamma t/2} e^{i\varphi} | 1 \rangle}{\sqrt{|\alpha|^2 + |\beta|^2 e^{-\Gamma t}}}, \text{ if not tunneled} \end{cases}$$

(better theory: Pryadko & A.K., 2007)

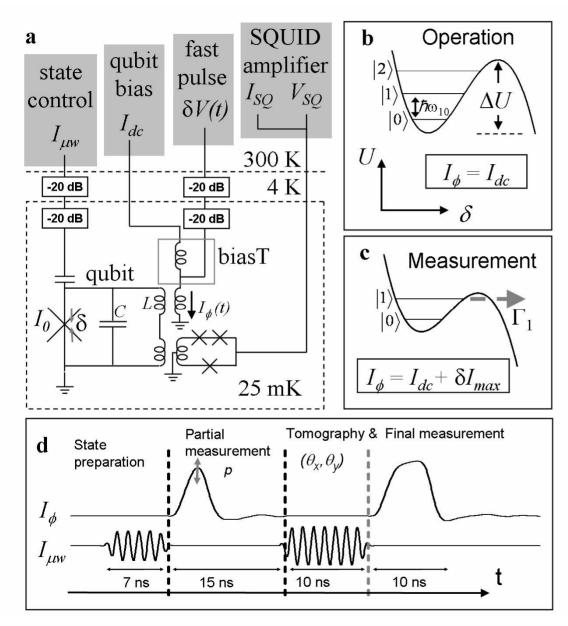
amplitude of state |0> grows without any physical interaction (!) continuous null-result collapse

(similar to optics (never realized) Dalibard-Castin-Molmer,1992)

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Experimental technique for partial collapse



Nadav Katz *et al*. (John Martinis' group)

Protocol: 1) State preparation (microwave pulse)

- 2) Partial measurement by lowering barrier for time t
- 3) State tomography (microwave + full measurement)

Measurement strength $p = 1 - \exp(-\Gamma t)$ (actually controlled by Γ)

p=0: no measurement
p=1: orthodox collapse
anything in between!

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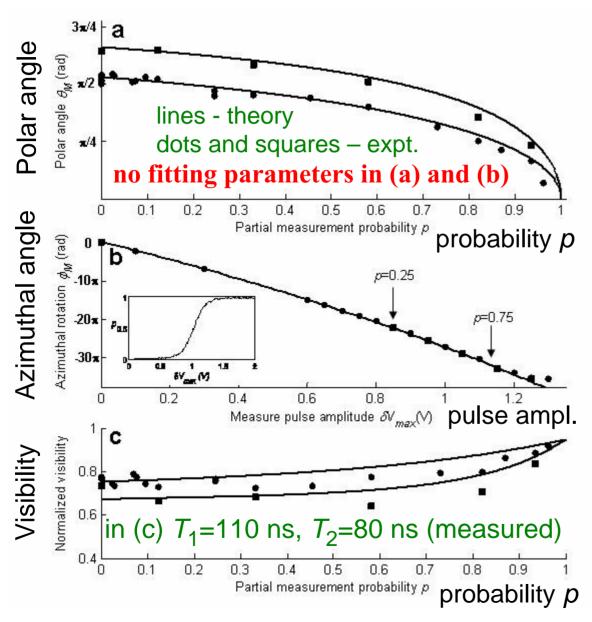
Experimental tomography data

Nadav Katz et al. (UCSB) Ψ_{in} p=0p = 0.14p = 0.06 $|0\rangle + |1\rangle$ No partial measure ٠θy -1 0 1 O quadrature amplitude [D₀1/x] -1 0 1 guadrature amplitude [0,.0x] -1 0 O quadrature amplitude [O,, I] p=0.43 p=0.32 p=0.23 θx $\frac{\pi}{\pi/2}$ -1 0 1 Q quadrature amplitude [D_101] -1 0 1 O quadrature amplitude [0, Uz] -1 0 1 O quadrature amplitude [D_Us] p = 0.83p=0.56 p = 0.701% pattial measurem -1 0 1 O quadrature amplitude [D_U/z] -1 0 1 O quadrature amplitude [D,,1/x] -1 0 1 O guadrature amplitude [0, 1/x]

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Partial collapse: experimental results



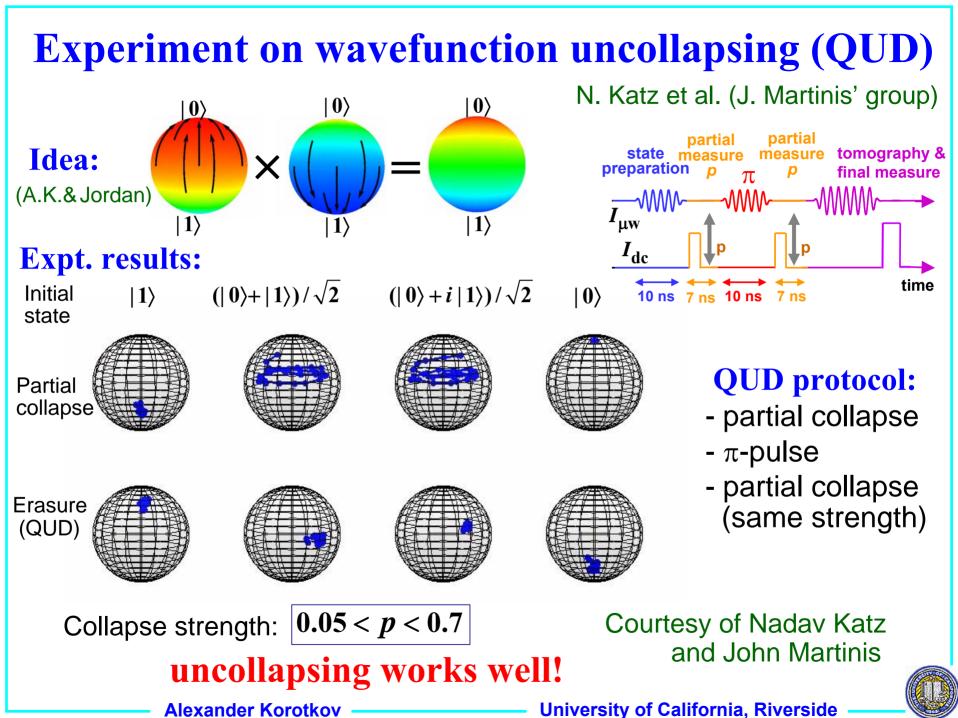
N. Katz et al., Science-06

• In case of no tunneling (null-result measurement) phase qubit evolves

• This evolution is well described by a simple Bayesian theory, without fitting parameters

• Phase qubit remains fully coherent in the process of continuous collapse (experimentally ~80% raw data, ~96% after account for T1 and T2)

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Conclusions

- Continuous (weak) quantum measurement is a new interesting subject in solid-state mesoscopics; interest grows
- A number of experimental predictions have been made, two direct experiments have been realized

Future directions

- Understanding of weak (continuous) quantum measurement for particular systems
- Analysis of possible applications (quantum feedback, realistic non-QND qubit measurements, non-unitary quantum gates, continuous error correction, etc.)
- Experimental proposals (realistic and for future)
- Actual experiments (direct and indirect; to resolve controversies and towards applications)



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