Two-qubit decoherence mechanisms revealed via quantum process tomography

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QPT basics

Quantum operation (completely positive linear map)

(i.e. positive Hermitian (d.m.) \rightarrow positive Hermitian (d.m.); "completely" means even with ancilla)

$$\rho = L[\rho^0], \quad \rho_{ij} = \sum_{k,l=0}^{d-1} L_{ij,kl} \rho_{kl}^0$$

Can be represented with Kraus operators (not uniquely): $L[\rho^0] = \sum_n K_n \rho^0 K_n^{\dagger}$ Unique representation: $L[\rho^0] = \sum_{m,n=0}^{d^2-1} \chi_{mn} E_m \rho^0 E_n^{\dagger}$ where E_m is a chosen (arbitrary) basis of operators "Pauli basis": products of (*I*,*X*, *Y*,*Z*) for each qubit (jargon: X= σ_X , etc.) Sometimes modified Pauli basis: (*I*,*X*, *-iY*,*Z*)

$$L[\rho^{0}] = \sum_{m,n=0}^{d^{2}-1} \chi_{mn} E_{m} \rho^{0} E_{n}^{\dagger}$$

- not quite intuitive
- Hermitian, $d^2 \times d^2$, positive-semidefinite
- Tr χ =1 (usually)
- for a unitary operation $U \quad \chi_{mn} = k_m k_n^*, \quad U = \sum_{n=1}^{d^2 1} k_n E_n$

Markovian decoherence (λ-matrix)

 $\frac{d}{dt}\rho = M\rho \quad \text{then } L = \exp(Mt) \quad (\text{if treat density matrix as a vector})$ $M = M_{\text{coh}} + D \quad \text{Can introduce } \lambda \text{-matrix:} \quad D[\rho] = \sum_{m,n=0}^{d^2-1} \lambda_{mn} E_m \rho E_n^{\dagger}$ $\lambda \text{-matrix is similar to } \chi \text{-matrix,}$ but Tr $\lambda = 0$ and different dimension ([χ]=[λ t])

If $M_{\text{coh}}=0$, then at weak decoherence $\chi \approx \chi^{I} + \lambda t$, $\chi_{mn}^{I} = \delta_{m0}\delta_{n0}$

Local vs. non-local decoherence

Uncoupled 2-qubit system and local decoherence: $\chi = \chi^{(1)} \otimes \chi^{(2)}$

Natural to introduce nonlocality parameter for decoherence:

$$\varepsilon_{NL} = \frac{\operatorname{Tr} |\chi - \tilde{\chi}|}{\operatorname{Tr} |\chi - \chi_{\text{ideal}}|}, \quad \tilde{\chi} = \tilde{\chi}^{(1)} \otimes \tilde{\chi}^{(2)}, \quad \tilde{\chi}^{(1)} = \operatorname{Tr}_{2} \chi \\ |A| \equiv \sqrt{A^{\dagger} A}, \quad \operatorname{Tr} |A| - \text{``trace norm''}$$

Not so easy for **coupled** system. However, in Markovian case can use λ -matrix:

$$\varepsilon_{NL}^{'} = \frac{\operatorname{Tr} |\lambda - \tilde{\lambda}|}{\operatorname{Tr} |\lambda|}, \quad \tilde{\lambda} = \tilde{\lambda}^{(1)} \otimes \chi^{I(2)} + \chi^{I(1)} \otimes \tilde{\lambda}^{(2)}, \quad \tilde{\lambda}^{(1)} = \operatorname{Tr}_{2} \lambda$$

 $\mathcal{E}_{NL} \approx \mathcal{E}_{NL}$ for no coherent evolution and weak decoherence

Considered models of decoherence

Energy relaxation

$$\begin{aligned} \lambda_{00} &= -2(1/T_1^{(1)} + 1/T_1^{(2)}) \\ \lambda_{11} &= \lambda_{22} = 1/T_1^{(2)}, \, \lambda_{44} = \lambda_{88} = 1/T_1^{(1)} \\ \lambda_{03} &= \lambda_{30} = 1/T_1^{(2)}, \, \lambda_{0,12} = \lambda_{12,0} = 1/T_1^{(1)} \\ \lambda_{21} &= -\lambda_{12} = i/T_1^{(2)}, \, \lambda_{84} = -\lambda_{48} = i/T_1^{(1)} \end{aligned}$$

Partially correlated (nonlocal) dephasing

$$\begin{split} \lambda_{00} &= -(1/T_{\varphi}^{(1)} + 1/T_{\varphi}^{(2)})/2 \\ \lambda_{33} &= 1/2T_{\varphi}^{(2)}, \, \lambda_{12,12} = 1/2T_{\varphi}^{(1)} \\ \lambda_{3,12} &= \lambda_{12,3} = -\lambda_{0,15} = -\lambda_{15,0} = \kappa/2\sqrt{T_{\varphi}^{(1)}T_{\varphi}^{(2)}} \\ \kappa \text{ - correlation factor } (\kappa=0 \text{ if local}) \end{split}$$

Noisy coupling

$$\begin{split} \lambda_{00} &= -\Gamma_{s} \\ \lambda_{55} &= \lambda_{10,10} = \lambda_{5,10} = \lambda_{10,5} = \\ &= \lambda_{0,15} = \lambda_{15,0} = \Gamma_{s} / 2 \end{split}$$

0,1,2, ...15 = *II*, *IX*, *IY*, *IZ*, *XI*, *XX*, *XY*, *XZ*, ... *ZZ*

Specific pattern for each model! (usually different elements)

If no evolution (memory), then easy:

$$\chi \approx \chi^{I} + \lambda t, \quad \chi^{I}_{mn} = \delta_{m0} \delta_{n0}$$

 $\lambda = \lambda_{enrel} + \lambda_{corrdeph} + \lambda_{noisy coupl}$

If coupled qubits, then should not be easy $\chi \neq \chi_{ideal} + \lambda t$

However, miraculously, for sqrt{iswap} and for largest extra elements of χ , it is still a simple addition, that makes distinguishing decoherence models very simple.



QPT for one-qubit uncollapsing

Kyle Keane (started recently)



Actually, problem with definition of χ -matrix, because probability depends on initial state